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ROTATIONALLY RESOLVED SPECTROSCOPY OF TETRACENE AND ITS VAN DER WAALS COMPLEXES WITH INERT GAS ATOMS

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Rovibronic spectra of tetracene ( $C_{18}H_{12}$ ) and its Van der Waals complexes with argon and krypton have been resolved by using a well collimated molecular beam in combination with a single frequency dye laser. The rotational constants of tetracene have been deduced in the  $S_0$  and  $S_1$  electronic state. The rovibronic spectra of the Van der Waals complexes show an enhanced intersystem crossing.

### 1. INTRODUCTION

During the last decade there has been an increasing interest in the study of large molecules and their van der Waals (VdW) complexes. Part of this interest is caused by the development of very sensitive laser-spectroscopic techniques in combination with effective reduction of the internal energy of such molecules in molecular beams. Most of these experiments use a free jet expansion in which the laser beam crosses the molecular beam very near the nozzle. For large aromatic molecules like anthracene or tetracene this allows vibrational resolution and provides a tool for the study of processes like intermolecular vibrational relaxation and

The rotational spectra of such large molecules can be resolved by strongly collimating the molecular beam, thus decreasing the residual Dopplerwidth. This provides a very high resolution and allows the study of for instance the fluorene-argon VdW complex<sup>2</sup> or even the revelation of the molecular eigenstates of pyrazine<sup>3</sup>.

Now we report the first observation of rotationally resolved spectra of tetracene (T) and its VdW complexes with argon and krypton. The T molecule and its VdW complexes have been studied under vibrational resolution  $^{1}$ ,  $^{4}$ ,  $^{5}$  and calculations have been reported on the structure of the T-Ar complex  $^{6}$ . An enhanced intersystem crossing in the complex has been concluded from lifetime measurements on the T-rare gas complexes  $^{4}$ . We affirm this perturbation of the  $S_{1}$  electronic state from our spectroscopic data of T-Ar and T-Kr. Besides we report the analysis of the rovibronic spectra of T from which we deduced molecular constants of this molecule in the  $S_{0}$  and  $S_{1}$  electronic state. All strong spectral features could be assigned and we found no evidence for perturbations of the  $S_{1}$  state like have been observed for pyrazine  $^{3}$  or are present in the T-rare gas complexes.

# 2. EXPERIMENTAL

In our studies we used the seeded beam technique, combining the advantage of a considerable reduction of the internal temperature of the molecules with a high density of VdW complexes. In molecules as large as T, such internal cooling is indispensable for reduction of the numerous spectral lines. The experimental set-up has been described in

detail before<sup>7,8</sup>. We used a quartz source in which a supply of T was heated to 210 °C, mixed with a noble gas and expanded in a vacuum chamber by a nozzle (100 µm diameter). The molecular beam was doubly skimmed to decrease the residual Doppler width to about 15 MHz in the interaction zone with the radiation field. The undispersed laser induced fluorescence was imaged to the photocathode of a PMT (EMI 9863/350 QA), monitored by a standard photon-counting system with an integration time of typically 0.1 sec.

We obtained the narrow band radiation field from a standard dye laser system (Coherent 599-21) with a 3 MHz bandwidth and a maximum output power of 50 mW. The bandwidth however could be increased to about 30 GHz by removing the intracavity etalons, thus allowing continuous laser scans (by tuning a Lyot filter) over the whole region of the gain medium.

For relative frequency marking we used a sealed off temperature stabilized Fabry-Perot interferometer with a 150 MHz free spectral range. Absolute frequency measurements have been performed with a home-built Michelson interferometer-type wavelength meter<sup>8</sup>.

# 3. RESULTS AND DISCUSSION

# 3.1. $\underline{S_0 \rightarrow S_1}$ transition in tetracene

The  $S_0 \rightarrow S_1$  electronic transition in tetracene was studied by Amirav et al under vibrational resolution<sup>1</sup>. Several vibrational levels were assigned in both the  $S_0$  and  $S_1$  state. At vibrational energies below  $E_v = 1000 \text{ cm}^{-1}$  isolated levels were observed while above this level Fermi resonances appear, fading into a quasi continuum at  $E_v = 1900 \text{ cm}^{-1}$ .

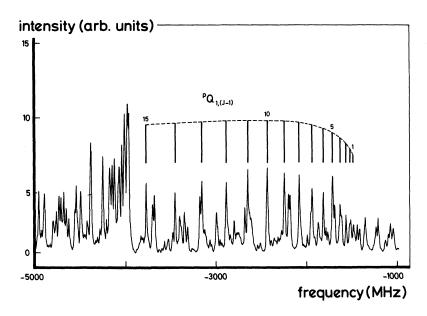


FIGURE 1 Part of the  $S_0 \rightarrow S_1$  rovibronic transition of tetracene at 22396.53(2) cm<sup>-1</sup>. The indicated lines belong to a Q-branch with  $\Delta k_{-1} = -1$ ,  $k_{-1} = 1$  and  $k_1 = J - 1$ .

We studied three vibronic transitions in the low energy range (below  $E_v = 1000 \text{ cm}^{-1}$ ) under rotational resolution (table 1). Two of these vibrations are totally symmetric  $a_g$ -type vibrational modes, which are active for the  $S_0$  ( $^1A_g$ )  $\rightarrow$   $S_1$  ( $^1B_{2u}$ ) symmetry allowed electronic transition. We also investigated a non-totally symmetric vibration of  $b_{3g}$  type. Transitions to this vibrational mode gain their intensity from a coupling with the higher energy  $^1B_{1u}$  electronic state and are of much less intensity as the symmetry allowed transitions.

The observed spectrum around 22397 cm<sup>-1</sup> was identified by Amirav et al. as the electronic origin of the  $S_0(^1A_g) \rightarrow$ 

TABLE 1 Excited state vibrational assignments and experimental transition frequencies of tetracene and the redshifts of transitions in the T-Ar and T-Kr VdW complexes

	T		T-Ar	T-Kr
ν <sub>0</sub> (cm <sup>-1</sup> )	assignment	symmetry	$\Delta v_0^{(cm^{-1})^a}$	Δν <sub>0</sub> (cm <sup>-1</sup> ) <sup>a)</sup>
22 396.53(2)	electronic origin	a g	41.67(5)	66.9(1)
22 707.84(2)	311 cm <sup>-1</sup>	a g	41.41(5)	66.5(1)
22 876.62(2)	471 cm <sup>-1</sup>	b <sub>3g</sub>	41.47(5)	

a) The accuracy is limited by the lack of a complete rotational assignment of the transitions.

 $S_1(^1B_{2u})$  transition. The rotational spectrum was identical to the one around 22 708 cm $^{-1}$ . They were uniquely identified as b-type perpendicular transitions, corresponding to a transition dipole moment parallel to the short molecular axis. The non-totally symmetric  $b_{3g}$ -type vibration though has a very different spectrum, which was identified as an a-type (parallel) transition. All the spectral features were independent of the type of seeding gas and therefore cannot be due to VdW complexes of T with inert gas atoms. The  $S_0 + S_1$  electronic origin has a very strong spectrum (fig. 1) and reduction of the laser power to as little as a few mW was necessary to avoid strong saturation effects.

We assigned a total of 230 lines in the various spectra (up to J=15) and fitted them to an asymmetric rotor model with an excellent fit. From the fit the rotational constants were obtained in the ground and excited state, the identification of the lines and the transition frequency  $v_0$ . In table 2 the rotational constants are listed for the ground

TABLE 2 Rotational constants of the tetracene molecule in the  $S_0$  and  $S_1$  electronic state (  $\Delta A = A' - A''$  ,  $\Delta B = B' - B''$  ,  $\Delta C = C' - C''$  ).

s <sub>0</sub>		s <sub>1</sub>
constant (MHz)		constant (MHz)
A	1630 (1)	ΔΑ 17.4 (1.2)
В	213.4(2)	$\Delta B = -1.81(8)$
С	188.8(2)	ΔC -1.19(8)

and excited electronic state. Within the experimental errors no differences in the rotational constants of the various vibrational states were observed.

The spectra of T are pure unperturbed rotational spectra and we could assign all relevant spectral features. This indicates that in the low vibrational energy range in the S, electronic state there is no clear intersystem crossing. The interaction between the S, state and background states is very weak although it is expected that a molecule as large as T should be in the "statistical limit"9. The lifetimes of the  $S_1$  states (20 ns)<sup>1</sup> are such that the expected natural linewidth is still smaller than our experimental linewidth. This is in accord with our observations.

# 3.1. Tetracene-argon VdW complexes

Tetracene-argon VdW complexes  $(T-Ar_n, n=1,2,3, ...)$  are easily formed in the molecular beam by using argon as seeding gas. Operating our laser in the broadband mode (30 GHz bandwidth) we observed vibrationally resolved spectra comparable with those of ref. 4 and ref. 5. The transitions in the com-

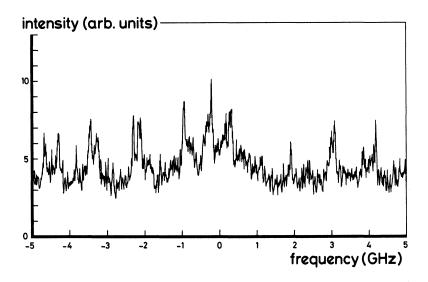


FIGURE 2 The  $S_0 \rightarrow S_1$  electronic origin of the tetraceneargon Van der Waals complex at 22 354.86(5) cm<sup>-1</sup>.

plexes are more or less regularly redshifted compared to the analogous transitions in T. Besides these shifts, the identification of the complexes is based on the pressure dependence of the intensity of the observed transitions.

Model calculations in the structure of the T-Ar complex<sup>6</sup> predict a favourable position of the Ar-atom above the molecular plane at a distance of about 3.45 Å and very near the middle of the T molecule. Although these calculations only give an indication of the structure of the complex we found similar calculations very satisfactory for the fluoreneargon VdW complex<sup>2</sup>. An experimental determination of the structure of this complex can be given by deriving the moments of inertia from the rotational spectra.

We recorded rotationally resolved spectra of transitions in T-Ar analogous to the vibronic transitions in T mentioned before (table 1). The spectrum around 22 354 cm<sup>-1</sup> is assigned as the electronic origin of the T-Ar complex (fig. 2). It is of perpendicular type in common with the electronic origin of T and it is almost identical to the transition at  $22.666 \,\mathrm{cm}^{-1}$ . The spectrum around  $22.826 \,\mathrm{cm}^{-1}$  belongs to a different vibrational mode of T and is of parallel type. The linewidth in these spectra is still determined by the residual Doppler width, in accordance with the expected natural linewidth of 6 MHz4. We didn't succeed in assigning the pronounced spectral features in the spectra of T-Ar. For instance the sequence of Q-branches in the perpendicular spectra looks distorted and there seems to be an excess of lines in these spectra. We are convinced that there are perturbations in the  $S_1$  electronic state due to couplings with background states. This is in accord with the lifetime measurements of ref. 4. The evidence for such coupling in these complexes is more clear in the T-Kr complex where the perturbations are much larger.

# 3.3. Tetracene-krypton VdW complexes

By using krypton as seeding gas we produced T-Kr $_{\rm n}$  (n = 1, 2, 3, ...) VdW complexes in the molecular beam. The transition frequencies (table 1) in the T-Kr complex have a different shift from the corresponding transitions in T than those of the T-Ar complex. This allows us to distinguish between the various molecules and is very convenient to check the identification of the transitions. The electronic origin of the T-Kr complex around 22 330 cm $^{-1}$  has a very striking spectrum (fig. 3). Only rotational contours were observed, due to the small rotational constants and a

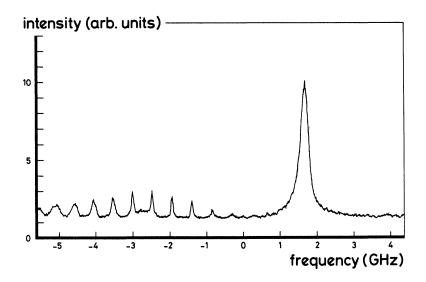


FIGURE 3 The  $S_0 \rightarrow S_1$  electronic origin of the tetracene-krypton Van der Waals complex at 22 329.6(7) cm<sup>-1</sup>.

presumably increased linewidth<sup>4</sup>. The spectrum contains a pronounced feature, resembling a Q-branch of a parallel transition, and a sequence of weaker branches, characteristic of a perpendicular transition. Besides this unusual contour there is a rather strong background in the spectrum.

The vibronic transition of T-Kr connected to the ag vibrational mode of T mentioned before is very different from the electronic origin of T-Kr. Note that this was not the case in the T-Ar or the T molecule. This spectrum is almost without any structure on a broad background. It has been demonstrated that the 311 cm vibrational excitations of T-Ar and T-Kr do not result in vibrational predissociation on the time scale of the excited state life-

times. We conclude that the perturbed spectra are due to an intersystem crossing which is almost absent in T but is enhanced by addition of the inert gas atoms.

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