Infrared Multiple Photon Dissociation Spectroscopy of a Gas-Phase Oxo-Molybdenum Complex with 1,2-Dithiolene Ligands

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* Supporting Information

ABSTRACT: Electrospray ionization (ESI) in the negative ion mode was used to create anionic, gas-phase oxo-molybdenum complexes with dithiolene ligands. By varying ESI and ion transfer conditions, both doubly and singly charged forms of the complex, with identical formulas, could be observed. Collision-induced dissociation (CID) of the dianion generated exclusively the monoanion, while fragmentation of the monoanion involved decomposition of the dithiolene ligands. The intrinsic structure of the monoanion and the dianion were determined by using wavelength-selective infrared multiple-photon dissociation (IRMPD) spectroscopy and density functional theory calculations. The IRMPD spectrum for the dianion exhibits absorptions that can be assigned to (ligand) C=C, C=S, C=C≡N, and Mo=O stretches. Comparison of the IRMPD spectrum to spectra predicted for various possible conformations allows assignment of a pseudo square pyramidal structure with $C_{2v}$ symmetry, equatorial coordination of MoO$_2^+$ by the S atoms of the dithiolene ligands, and a singlet spin state. A single absorption was observed for the oxidized complex. When the same scaling factor employed for the dianion is used for the oxidized version, theoretical spectra suggest that the absorption is the Mo=O stretch for a distorted square pyramidal structure and doublet spin state. A predicted change in conformation upon oxidation of the dianion is consistent with a proposed bonding scheme for the bent-metalloocene dithiolene compounds [Lauher, J. W.; Hoffmann, R. J. Am. Chem. Soc. 1976, 98, 1729–1742], where a large folding of the dithiolene moiety along the S···S vector is dependent on the occupancy of the in-plane metal d-orbital.

INTRODUCTION

Among the mononuclear molybdenum enzymes, the DMSO reductase family is the most diverse in terms of their structure and function. Members of this family are involved in global C, S, N, and As cycling; for example, formate dehydrogenase catalyzes the transformation of formate to CO$_2$, DMSO reductase catalyzes the reduction of dimethyl sulfoxide (DMSO) to dimethyl sulfi(de (DMS), nitrate reductase catalyzes the reduction of nitrate to nitrite, and arsenite oxidase catalyzes the oxidation of arsenite to arsenate.$^{1-4}$ These substrate transformations require transfer of two electrons, and during catalysis the Mo center shuttles between the +4 and +6 oxidation states. The catalytically competent Mo center is regenerated by two one-electron steps, thereby transiently passing though the +5 state. The different states of the metal center have been spectroscopically (e.g., EXAFS, EPR, and resonance Raman) characterized.$^{5-9}$ In all cases, in the fully oxidized state, the Mo center is coordinated by two ene-dithiolate moieties from pyranopterin cofactors (Figure 1). In most cases, a terminal oxo group occupies the fifth coordination site, and depending on the enzyme, the Mo center is coordinated by additional ligands such as serine (in DMSO reductase), cysteine (in nitrate reductase), selenocysteine (in formate dehydrogenase), or hydroxide (in arsenite oxidase).
These coordination modes have been confirmed by protein crystallography.10,11

The spectroscopic and structural studies to date have provided the motivation to develop synthetic molecules with the \{Mo(O)S₄\} core as a vehicle to understand the intricate details of the electronic structure of the enzymatic Mo center.12–15 The terminal oxo coordination provides a strong ligand field that orients the redox active dₓₒ−yₒ orbital in the equatorial plane, which implies that the ene-dithiolene moieties are involved in the electron transfer pathway to and from the Mo center. Mo compounds possessing a paramagnetic \{Mo²⁺O₄S₄\} core have been used to probe \(S → Mo\) charge transfer transitions by a variety of spectroscopic approaches such as EPR, absorption, MCD, and resonance Raman spectroscopy.

Early work by Lauher and Hoffmann provided a bonding scheme for bent-metallocene dithiolene compounds,16 where a large folding of the dithiolene moiety along the S−S vector has been described. They proposed that the magnitude of the folding is dependent on the occupancy of the in-plane metal d-orbital; i.e., the more occupied the orbital is, the less folding is observed (Figure 2). Enemark and co-workers described the folding of the dithiolene ligand in o xo-Mo centers as a contributing factor to modulation of electron transfer in mononuclear molybdenum enzymes during the regeneration of the catalytically competent state.17–19 In this case, the redox orbital interacts with the symmetric combination of out-of-plane orbitals of the dithiolene sulfur.

The intrinsic propensity of folding of a dithiolene ligand as a function of the oxidation state of the metal can be measured best in discrete model o xo–metal–ligand complexes, without any influences from counterion, solvent, or other condensed-phase effects. Vibrational spectroscopy in the gas phase provides a means to study the intrinsic structure and behavior of metal–ligand complexes, and it is now well-known that infrared spectra of ionic species confined to the gas-phase environment of a mass spectrometer can be collected using the combination of tandem mass spectrometry and (wavelength-selective) infrared multiple-photon dissociation (IRMPD) spectroscopy [for reviews, see refs 20–26 and the literature cited therein]. Much of the work involving IRMPD spectroscopy and metal complexes has focused on probing the interaction(s) between cations and biologically relevant molecules.27–30 In many cases, the attention has been on the intrinsic structure and bonding interactions of inorganic and organometallic species.31–64

In this study we used electrospray ionization (ESI) to generate discrete, gas-phase anions from a model oxo-molybdenum complex with dithiolene ligands with formula \([MoO(mnt)₂]^{2−}\), where \(n = 1\) or 2 (mnt\(^{2−}\) = 1,2-dicyanoethylenedithiolate; Figure 2). The ions were then investigated in the gas phase using collision-induced dissociation (CID) and IRMPD. In the IRMPD spectroscopy experiment used in this work, gas-phase ions are irradiated at mid-IR wavelengths using a free electron laser (FEL). When the FEL wavelength matches the energy of a vibrational transition, absorption of multiple photons raises the vibrational energy of the trapped ion to the dissociation threshold. An IR spectrum is generated by measuring the fragmentation induced by this process as a function of photon wavelength. Vibrational mode assignment and structural determination is made with the assistance of frequencies predicted by density functional theory (DFT) or related computational methods.

Our primary goals in this study were (a) to use wavelength-selective infrared multiple-photon photo dissociation spectroscopy to determine the structure of the dianionic complex \([MoO³⁻O(mnt)₂]^{2−}\) and (b) measure the intrinsic frequency of the Mo=O stretch within the complex. The results obtained will pave the way for more detailed investigations of the nature of metal–ligand interactions using a library of similar complexes designed to mimic the important oxo-transfer chemistry of the pyranopterin enzymes.

### EXPERIMENTAL METHODS

**Collision-Induced-Dissociation (CID) Experiments.** The tetraethylammonium salt of the molybdenum(IV) complex \([Et₄N][MoO²⁻O(mnt)₂]\) was synthesized using an established procedure described in detail elsewhere.65 Preliminary ESI and CID experiments were performed on a ThermoScientific (San Jose, CA) LTQ-XL linear ion trap mass spectrometer (MS) equipped with an Ion Max ESI source. For the ESI experiments, a stock solution (approximately 0.001 M) of the \([MoO³⁻O(mnt)₂]^{2−}\) complex as the tetraethylammonium salt was prepared in acetonitrile. The solution was infused into the ESI-MS instrument using the incorporated syringe pump at a flow rate of 10–15 μL/min. The atmospheric pressure ionization stack settings for the LTQ (lens voltages, quadrupole and octopole voltage offsets, etc.) were optimized for maximum transmission of singly or doubly charged anions to the ion trap mass analyzer by using the autotune routine within the LTQ Tune program. In general, harsher ESI and ion transmission conditions (i.e., higher tube lens and skimmer voltages) result in higher yields of the singly charged ion, which is presumably created by electron detachment from the dianion. Helium was used as the bath/ buffer gas to improve trapping efficiency and as the collision gas for CID experiments.

Because of the complex isotopic pattern of Mo, precursor ions were isolated for the initial dissociation stage (MS/MS) using an isolation width of 0.9–1.2 mass to charge (m/z) units centered on the \(^{98}\)Mo isotopic peak. Product ions selected for subsequent CID (MS³ experiments) were isolated using slightly greater widths (1.2–1.5 m/z units) to improve trapping and fragmentation efficiency. For each stage, the exact width was chosen empirically to produce the best compromise between high precursor ion intensity and the ability to isolate a single
isotopic peak. The (mass) normalized collision energy (NCE, as defined by ThermoScientific) was set between 25 and 35%, which corresponds to the application of roughly 0.55–0.68 V tickle voltage to the end-cap electrodes with the current instrument calibration. The activation Q, which defines the frequency of the applied rf potential, was set at 0.30. In all cases, the activation time employed was 30 ms. Spectra displayed represent the accumulation and averaging of at least 30 isolation, dissociation, and ejection/detection steps.

**ESI FT-ICR Mass Spectrometry.** As in the preliminary ESI studies, a stock solution (approximately 1.0 × 10^{-4} M) of the \([\text{MoO}(\text{mnt})_2]^{2-}\) complex was prepared in acetonitrile for the IRMPD experiments. Previously established methods used by our group for generation of ions and the subsequent collection of IRMPD spectra^{58–64} were used here. Briefly, ESI was performed using a Micromass (now a component of Waters Corp., Milford, MA) Z-Spray source in the negative ion mode. Dry nitrogen (∼80 °C) was used to assist in the desolvation process. Ions were injected into a home-built Fourier transform ion cyclotron resonance (FT-ICR) mass spectrometer described in detail elsewhere.\(^{66}\) Ions were accumulated for the duration of the previous FT-ICR cycle (6 s) in an external hexapole and injected into the ICR cell via a quadrupole deflector and an octapole rf ion guide. Instrument operating parameters, such as desolvation temperature, cone voltage, and ion accumulation and transfer optics voltages, were optimized to maximize the formation of either singly or doubly charged anions and transfer of the species to the ICR cell.

**Infrared Multiple Photon Dissociation (IRMPD).** Infrared spectra were recorded by measuring the photodissociation yield as a function of photon wavelength. Precursor anions were irradiated using 23 FELIX macropulses (35 mJ/macropulse, 5 μs pulse duration, full width at half-maximum (fwhm) bandwidth ~0.5% of central λ). In the IRMPD process, a photon is absorbed when the laser frequency matches a vibrational mode of the gas-phase ion and its energy is subsequently distributed over all vibrational modes by intramolecular vibrational redistribution (IVR). The IVR process allows the energy of each photon to be dissipated before the ion absorbs another, which leads to promotion of ion internal energy toward the dissociation threshold via multiple photon absorption.\(^{67}\) It is important to note that infrared spectra obtained using IRMPD are comparable to those collected using linear absorption techniques.\(^{68,69}\) For these experiments, the FEL wavelength was tuned between 5.7 and 14 μm in 0.03–0.1 μm increments. The intensities of product and undissociated precursor ions were obtained from an averaged mass spectrum measured using the excite/detect sequence of the FT-ICR-MS after each IRMPD step. The IRMPD yield was normalized to the total ion current.

**DFT Geometry and Frequency Calculations.** All DFT calculations were performed using the Gaussian 09 program.\(^{70}\) Initial optimization of \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\) and \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\) was performed at the B3LYP/3-21G* level of theory using geometries in which dithiolene ligands were randomly arranged around a \{MoO^{**}\} core and coordination of the metal was by either the thiolate or cyano groups. Minima identified after the initial calculations were then subjected to full optimization using the same functional, effective core potential, and associated basis on Mo (MWB28) and the 6-311+G(d) basis set on C, N, O, and S.

To test the general consistency of relative rankings of energy for various minima, calculations were also performed with the 6-311+G(3df) basis set on C, N, O, and S. An exhaustive survey of models, functionals, and basis sets is beyond the scope of this investigation. However, to check general agreement between experiment and theory, bond lengths and vibrational frequencies were also calculated using the M06-L functional. The hybrid B3LYP functional, a standard in IRMPD investigations of gas-phase ions, is an approximation to the exchange–correlation energy functional which includes some portion of exact exchange from Hartree–Fock theory with exchange and correlation from other sources.\(^{71–74}\) M06-L is one of a group of meta-GGA functionals.\(^{75,76}\) We have found in prior studies of gas-phase metal complex thermochemistry that the M06 functional accurately reproduced trends in ligand exchange and addition to uranyl species.\(^{77}\) The M06-L is fully local, with no Hartree–Fock exchange, and is also reported to be effective for metal ions and inorganic and organometallic species.\(^{75}\)

The DFT calculations performed here were primarily to assist with assignment of vibrational modes and determinations of intrinsic structures. Our intent was not to model fragmentation energetics or rigorously determine the differences in energy of possible spin states of the \([\text{MoO}(\text{mnt})_2]^{1+}\) complexes. Therefore, no corrections were made for possible errors in the energies for the respective species due to differences in spin–orbit coupling for \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\) and \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\). In any case, any error that may be due to differences in spin–orbit coupling is expected to be minor (on the order of 200–600 cm^{-1}), particularly when compared to the overall electronic energy differences between \([\text{Mo}^{5+}\text{O}-(\text{mnt})_2]^{2-}\) and \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\) or difference spin states for the latter species, and should not hinder the qualitative determination of the preference for a given geometry or spin state.

Scaling factors were chosen empirically by bringing the predicted stretching frequencies in the C=C stretch region (ca. 1460 cm^{-1}) into agreement with the IRMPD spectrum. For a given functional/basis set combination, the same scaling factor, 0.96 at the B3LYP/6-311+G(d) level, was used for the spectra predicted for \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\) and \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\).

## RESULTS AND DISCUSSION

**ESI and CID of \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\) and \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\).** ESI mass spectrometry has been employed to investigate oxomolybdenum complexes with or without dithiolene ligands. For example, Dessapt and co-workers studied the formation of novel Mo(V) dithiolene compounds created by adding alkynes to solutions of MoO_S_2^{2-} in a mixture of MeOH and NH_3.\(^{79}\) Llusar and co-workers reported a combined ESI mass spectrometry and DFT study of sulfur-based reactions in MoS_4^{2-} and MoS_4^{2-} clusters that included dithiolene ligands.\(^{80}\) The structure of a phosphine oxide bound intermediate molecule originating from a dioxo-molybdenum-(VI) complex was investigated by ESI mass spectrometry and surface induced dissociation (SID).\(^{81}\) In addition, ESI mass spectrometry has been used in probing dynamics in oxygen atom transfer reactions.\(^{82,83}\) Important to our study is the fact that coupling of an ESI source to a photoelectron spectrometer\(^{84}\) has allowed a number of negatively charged inorganic species including a series of gas-phase oxomolybdenum(V)-tetrathiolate and bis(dithiolene) anions,\(^{85}\) and \([\text{Mo}^{5+}\text{O}(\text{mnt})_2]^{2-}\) (M = Mo and W)\(^{86}\) to be investigated by photodetachment photoelectron spectroscopy.
ESI mass spectra generated in our study from \([\text{Mo}^{IV}\text{O}-(\text{mnt})_2]^{2-}\) dissolved in acetonitrile are shown in Figure 3. In Figure 3a, the ESI and ion transfer conditions/voltages were tuned with the LTQ Tune program to maximize the intensity of the doubly charged anion \((m/z\ 194−199)\), which was observed as the most abundant ion. Also observed was a species that corresponds to a tetraethylammonium adduct to the dianion at \(m/z\ 518−528\). The inset in Figure 3a shows an expanded view of the high-resolution scan in the region of \(m/z\ 194−199\), which reveals the isotopic distribution expected for Mo, and a peak spacing of 0.5 mass unit to confirm the assignment of the \(-2\) charge state to the anion.

The spectrum in Figure 3b was generated by instead tuning the ESI and ion transfer voltages to optimize the yield of the singly charged anion (ca. \(m/z\ 394\)). Under these conditions, the singly charged anion was the dominant species in the ESI spectrum. Minor peaks (less than 10% relative intensity) were observed in the region of \(m/z\ 125−200\) which may correspond to Mo with a mixture of oxygen and sulfur atoms. The resolving power and mass measurement accuracy is not sufficient to distinguish oxygen from sulfur coordination in these ions. In Figure 3b, other minor peaks in the region \(m/z\ 200−350\) correspond to fragments of the monoanion. The tetraethylammonium adduct of the dianion \((m/z\ 518−528)\) and a lower mass species at \(m/z\ 417−427\) were also observed. CID (MS/MS stage, Figure S1a in the Supporting Information) of the species at \(m/z\ 524\) generates the ion at \(m/z\ 423\) through a neutral loss of 101 mass units (u). The product ion at \(m/z\ 423\) in turn is 29 u higher in mass than the species at \(m/z\ 394\). The loss of 101 u from the precursor at \(m/z\ 524\) is consistent with elimination of triethyl amine, and transfer of an ethyl group to the MoO-dithiolene anion. Subsequent CID of the product ion at \(m/z\ 423\) (MS\(^3\) stage, Figure S1b in the Supporting Information) furnishes the ion at \(m/z\ 394\). It is not clear from the MS/MS experiments whether the ethyl group is transferred to the Mo metal center, or instead to a dithiolene ligand.
Our primary interest in this study were the doubly and singly charged forms of the MoO-dithiolene complex. CID (MS/MS stage, Figure S2a in the Supporting Information) of the peak at m/z 197 produces exclusively the singly charged anion at m/z 394, presumably through electron detachment from one of the dithiolene ligands. Similar behavior of similar doubly charged

Figure 4. Predicted structures for [MoO(mnt)2]^{2−} and [MoO(mnt)2]^{−}. Spin states and energies for the species are provided in Table 1. Energies shown in the figure are relative to structure I and are the result of calculations at the B3LYP/6-311+G(d) level of theory.

Table 1. Electronic Energies for [Mo(mnt)2]^{2−} and [Mo(mnt)2]^{−}

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<th>ZPE</th>
<th>E + ZPE</th>
<th>ΔE (kcal/mol)</th>
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The ΔE values are relative to the lowest energy structure for [Mo(mnt)2]^{2−}.

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anions containing Mo and dithiolene ligands was reported by Llusar and co-workers. Subsequent CID (MS3 stage, Figure S2b in the Supporting Information) of the product ion at m/z 394 generated fragment ions at m/z 308, 312, 318, 330, and 366. The peak at m/z 366 is derived from a neutral loss of 28 u, which is attributed to elimination of CO by a mechanism that, while not clear, presumably involves decomposition of a dithiolene ligand and reaction with the oxygen atom of the Mo center. The loss of 64 u to generate the product ion at m/z 330 likely involves elimination of two S atoms. Likewise the fragment ion at m/z 318 may be generated by elimination of S and a C atom from the dithiolene ligand. The neutral losses associated with generation of the fragment ions at m/z 308 and 312 are consistent with elimination of S, CN, and CO (86 u) and S, C2, and CN (82 u), respectively. The apparent decomposition of the dithiolene ligands during CID warrants further study, and determination of the fragmentation pathways and their mechanisms would require isotope labeling.

DFT Calculations of Ion Structure. Potential structures for [MoIVO(mnt)2]2− and [MoIVO(mnt)2]− are shown in Figure 4. Relative energies for the respective species are provided in Table 1. Initial optimization of [MoIVO(mnt)2]2− was done using a singlet spin state. Three minima were identified, structures I, III, and IV, which differed in coordination of MoIVO by the dithiolene ligands. Structures in which the ligands coordinate the metal center through the cyanato groups (structures III and IV), were investigated for the sake of completeness and were found to be ~49–105 kcal/mol higher in energy at the B3LYP/6-311+G(d) level of theory than structure I (S coordination). Because the complexes with coordination by the cyanato groups are unlikely for the mnt ligands and energetically noncompetitive, and the predicted IR spectra are in poor agreement with the IRMPD results described below, further investigation using other functional and basis set combinations was not pursued.

Regardless of the functional or basis set used, the calculations performed here predict that the lowest energy structure for [MoIVO(mnt)2]2− is a pseudo square pyramidal conformation in which the oxo ligand occupies an axial position. The Mo atom lies slightly above the plane defined by the four (equatorial) S atoms of the dithiolene ligands. The experimental structural parameters for [MoIVO(mnt)2]2− from the Cambridge Structural Database, and the structural parameters computed in this work for both [MoIVO(mnt)2]2− and [MoIVO(mnt)2]−, are given in Table 2. Dihedral angles for the square pyramidal structure, measured using O axial−Mo−S ligand−C ligand, are provided in Table S1 of the Supporting Information. The experimental Mo≡O distances for [MoIVO(mnt)2]2− vary from 1.669 to 1.714 Å, a range of 0.045 Å. The calculated distance of 1.701 Å matches well with the experimental data, with the maximum deviation from the experimental values being 0.032 Å. As expected, in the Mo(V) complex the calculated Mo≡O distance shortens by 0.007 Å. Because of the symmetry all four Mo–S distances in [MoIVO(mnt)2]2− are equal. In the crystallographically determined structure of [MoIVO(mnt)2]2− the Mo–S distances vary from 2.370 to 2.417 Å, a range of 0.047 Å. The calculated Mo–S distance in [MoIVO(mnt)2]− is 2.443 Å, which is slightly longer than any of the experimental values with a maximum deviation from the experimental value of 0.07 Å. In the case of calculated C–S and C–C lengths for [MoIVO(mnt)2]2−, the maximum deviations from the experimental values are 0.086 and 0.060 Å, respectively, while the maximum deviation among the experimental values is 0.090 Å in both cases. For the Mo≡O angles in [MoIVO(mnt)2]2−, the experimental values range from 10.36 to 17.34°, a difference of 6.98°, while the maximum difference between the computational and experimental values is 4.49°. Optimized structures are dependent on many factors, such as the basis set and methodology used, e.g., ab initio, semiempirical, DFT, or molecular mechanics, while experimental structures may exhibit condensed-phase effects such as packing forces that result in deviations from gas-phase structures. While there are no well-defined criteria for comparing the “quality” of computed geometries, the differences in bond lengths and angles between the calculated and experimentally determined structures are similar to the variations observed among the experimental structures. We have discussed this issue in detail elsewhere.

A search was also conducted for a minimum corresponding to [MoIVO(mnt)2]2− in a triplet spin state. Structure II was found using both the B3LYP and M06-L functionals. A change in dihedral angles to 108.6° and 99.24° reflects rotation of the dithiolene ligands such that a structure more reminiscent of a distorted trigonal bipyramid is created. DFT predicts the triplet state to be ~32–34 kcal/mol higher in energy relative to the singlet state regardless of the functional or basis set used. However, we found that the triplet structure has one imaginary frequency, regardless of the level of theory or convergence criterion used, which indicates that it represents a saddle point rather than a true minimum. On the basis of this observation, we conclude that the most probable gas-phase conformation for [MoIVO(mnt)2]2− is pseudo square pyramidal, which is consistent with the local coordination environment in the Mo enzymes such as sulfite oxidase and members of the DMSO.
reductase family in catalytically competent Mo(IV) and Mo(V) states as revealed by X-ray absorption and EPR spectroscopies and X-ray crystal structures.5,6,11,89,90

Our results are in accord with more comprehensive calculations that were part of a photodetachment photoelectron spectroscopy study of the [MoIVO(mnt)2]2− by Wang, Wedd, and co-workers,86 in which the geometry of [MoIVO(mnt)2]2− was optimized under constrained C2v point symmetry. Using their calculations, the highest occupied molecular orbitals (HOMOs) of [MoIVO(mnt)2]2− were predicted to be of a1 symmetry and based upon the metal d_x^2−y^2 orbital. The next six molecular orbitals are different linear combinations of ligand sulfur 3p atomic orbitals that are symmetry-adapted for interaction with M and O atomic orbitals, some also with contributions from C 2p atomic orbitals involved in the C=S bond of the dithiolene unit (π or π*).86

Using conventional dithiolene coordination in our calculations, optimized structures for doublet and quartet spin states were identified for the singly charged anion (structures V and VI, respectively, in Figure 4). The lowest energy structure corresponds to the doublet, which features a square pyramidal conformation similar to the doubly charged anion except for a change in the overall orientation of one dithiolene ligand, which results in conversion from C2v to C3v symmetry. This is in accord with the previous photoelectron spectroscopy study,86 in which the oxidized complex was modeled in a doublet state (considered the likely configuration based on spin selection rule, ΔS = ±1/2). The calculations suggested that the single unpaired electron in [MoIVO(mnt)2]2− occupies the a1, d_x^2−y^2-based molecular orbital that corresponds to the doubly occupied HOMO of the parent dianions.

In our study, we also considered the quartet configuration for [MoIVO(mnt)2]−. As for the doubly charged anion, increasing the spin state of the singly charged anion resulted in a shift to a distorted trigonal-bipyramidal structure. Our calculations suggest that the quartet spin state is 43–45 kcal/mol less favorable than the doublet state when compared to the lowest energy structure for the doubly charged anion. As noted earlier, we assume that the differences in energy between the mono- and dianionic species are significantly greater than any error associated with differences in spin–orbit coupling.

Unlike [MoIVO(mnt)2]2−, no crystal structure data exists for [MoIVO(mnt)2]−, thus preventing an assessment of the accuracy of the calculated structures in terms of bond lengths and angles and folding angle. The DFT calculations predict a change in fold angle with oxidation. As noted earlier, Lauher and Hoffmann presented a bonding scheme for bent-metalloocene dithiolene compounds,16 and Enemark and co-workers later suggested that large folding of the dithiolene moiety along the S···S vector may be dependent on the occupancy of the in-plane metal d-orbital. The change in conformation predicted by our calculations likely represents the same type of ligand folding and is in accord with a change in conformation predicted in the photoelectron spectroscopy study,86 in which

Figure 5. (a) Experimental IRMPD and (b) predicted spectra for [MoO(mnt)2]2−. Theoretical frequencies were generated at B3LYP/6-311+G(d) level of theory using a singlet spin state and are scaled by a factor of 0.96.
The geometries with differentially folded dithiolene ligands were found to be favored by approximately 2 kcal/mol over the \( C_2v \) structure. The reduction from \( C_2v \) to \( C_s \) point symmetry apparently allows significant covalent interaction between the vacant metal-based \( d_{x^2-y^2} \) orbital and the highest energy \( S(\pi) \) orbital. The same orbitals cannot mix under \( C_2v \) point symmetry.

**IRMPD Spectroscopy.** The IRMPD spectrum for \([\text{MoO}(\text{mnt})_2]^{2-}\) was generated by monitoring formation of \([\text{MoO}(\text{mnt})_2]\) from electron detachment of \([\text{MoIVO}(\text{mnt})_2]^{2-}\) as a function of IR photon frequency and is shown in Figure 5a. The vibrational spectrum predicted for structure I, \([\text{MoIVO}(\text{mnt})_2]^{2-}\) with \( S \) coordination in the singlet state, is indicated by the dark trace in Figure 5b with the IRMPD spectrum included in gray to facilitate comparison. The spectrum obtained using the B3LYP functional and the 6-311+G(d) basis set on S, C, O, and N is used in Figure 5b because of the good qualitative and quantitative agreement with the IRMPD spectrum. Results obtained using the M06-L functional, and calculations at the B3LYP/6-311+G(3df) level of theory were in reasonable agreement with the experiment and are included in the Supporting Information.

As noted earlier, the structure identified for the triplet state appears to be a saddle point rather than a true minimum, so consideration of its vibrational features here is not appropriate. Comparisons of the IRMPD spectrum of \([\text{MoIVO}(\text{mnt})_2]^{2-}\) to spectra predicted for structures III and IV, those with coordination by one or both dithiolene ligands through the

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**Figure 6.** (a) Experimental IRMPD and (b and c) predicted spectra for \([\text{MoO}(\text{mnt})_2]^{2-}\). Theoretical frequencies were generated at the B3LYP/6-311+G(d) level of theory and are scaled by a factor of 0.96. The predicted spectrum in (b) is for a doublet spin state, while the one in (c) is for a quartet spin state.
cyano groups, are provided in Figure S3 of the Supporting Information. As noted earlier, structures III and IV are inconsistent with spectroscopic and crystallographic experiments in the condensed phase and are predicted to be significantly higher in energy by DFT.

Because of experimental limitations, the laser frequency could not be scanned into the IR region of the –C≡N stretch. However, three strong absorptions are observed in the IRMPD spectrum for [MoO\(_5\)O(mnt)$_2$]$^{2-}$ at 941, 1093, and 1464 cm$^{-1}$. Based on inspection of the vibrational modes predicted by DFT, these absorptions can be assigned to the Mo=O stretch, antisymmetric C–C≡N stretch, and antisymmetric (between the two ligands) C≡C stretch, respectively.

Two additional vibrations for [MoO\(_5\)O(mnt)$_2$]$^{2-}$ were predicted by DFT to appear at 838.5 and 966.9 cm$^{-1}$, and these modes correspond to antisymmetric and symmetric C=S stretches, respectively. It is clear from the predicted spectrum shown in Figure 5b that both C=S stretches are of low intensity. The reason for the absence of these peaks in the (experimental) IRMPD spectrum is not clear. We note that, while DFT in general may predict accurately the positions of absorptions in the IRMPD experiment, the predictions of intensity are less accurate. This observation may be attributed to the approximate nature of the harmonic frequency calculations. In addition, an IRMPD action spectrum is not identical to a linear absorption spectrum. While it is assumed that a linear proportionality between IRMPD yield and IR absorption intensity is a useful approximation, caution should be exercised in interpreting IRMPD spectra by comparison with calculated linear absorption spectra. These effects may include general red shifts and broadening of bands and changes in the relative intensities of bands due to the presence of nearby bands. Both general effects may be the result of anharmonicities of the vibrational modes in combination with the large number of photons (i.e., typically tens to hundreds) that are absorbed in the IRMPD process.

The IRMPD spectrum of [MoO\(_5\)O(mnt)$_2$]$^{-}$, shown in Figure 6a, was collected by monitoring the loss of 28 and 76 u from [MoO\(_5\)O(mnt)$_2$]$^{-}$ as a function of IR frequency (these fragmentation channels are consistent with those observed using CID). The spectrum for [MoO\(_5\)O(mnt)$_2$]$^{-}$ contains only a single peak at 960.6 cm$^{-1}$. Using the same scaling factor employed for the spectra predicted for [MoO\(_5\)O(mnt)$_2$]$^{2-}$, the predicted frequencies for [MoO\(_5\)O(mnt)$_2$]$^{-}$ are plotted in parts b and c of Figure 6 for the structures in the doublet and quartet spin states, respectively. With the scaling factor employed, there is good agreement between the position of the single absorption in the IRMPD spectrum and the position predicted for the Mo=O stretch in the oxidized complex with doublet spin state.

One explanation for the single absorption apparent in the spectrum of [MoO\(_5\)O(mnt)$_2$]$^{-}$ is that the barriers for photodissociation may be sufficiently high that they are only accessible via the intense Mo=O stretch. A similar explanation was used in an earlier study of the photodissociation of strongly bound uranyl anion complexes. In particular, the spectrum generated by loss of OCH$_3$ from uranyl methoxide contained only the antisymmetric uranyl peak, presumably because the energetics for the neutral loss were high and only accessed via the high-intensity $v_3$ uranyl absorption. We note that, in the present study, significantly higher relative collision energies were required to fragment the monoanion compared to the dianion, even when factoring in the difference in charge state. This would suggest that the barrier height to dissociation may play a role in determining the features of the IRMPD spectrum for [MoO\(_5\)O(mnt)$_2$]$^{-}$.

With a single peak in the IRMPD experiment, it is not possible to unambiguously assign the structure for [MoO\(_5\)O(mnt)$_2$]$^{-}$. However, based on the general intensities of absorptions predicted for the quartet spin state (with the caveat about predicted intensities noted above), one would expect more features to have appeared in the IRMPD spectrum if the quartet had been generated. More importantly, a shift of the Mo=O stretching frequency of 20 cm$^{-1}$ is predicted if one assumes oxidation to result in formation of the doublet state, in excellent agreement with the observed shift in the frequency in the experimental spectra. The shift to the blue by 20 cm$^{-1}$ upon oxidation of the complex can then be rationalized by the decrease in donation by a dithiolene group to the metal center, thus strengthening the Mo=O bond.

**CONCLUSIONS**

To summarize, ESI in the negative ion mode was used to generate gas-phase doubly and singly charged anions of an oxomolybdenum(IV) complex with dithiolene ligands. CID of the dianion generated exclusively the monoanion by electron detachment. Dissociation of the monoanion occurred through multiple pathways that include ligand elimination and ligand decomposition. Attempts were made to determine the structures of both the dianion and monoanion using IRMPD spectroscopy and DFT calculations. The photodissociation experiments, with comparison to predicted vibrational patterns, strongly suggest that the gas-phase structure of the dianion is a distorted square planar configuration, with the oxo ligand occupying an axial position and equatorial coordination of Mo by the dithiolene ligands through sulfur atoms. The particular conformation is in excellent agreement with the structures of other model compounds for DMSO reductase (and related enzymes) active sites and molybdenum enzymes as revealed by spectroscopic and crystallography studies. One peak was observed in the IRMPD spectrum for the oxidized, monoanionic species, thus preventing an unequivocal assignment of structure to either the doublet or quartet state structures predicted for the ion. However, the single absorption may reflect the fact that the majority of vibrational modes predicted for the lower-energy doublet state are low intensity, and the barrier for photodissociation reactions for the ion may be high. In addition, based on the general intensities of absorptions predicted for the quartet spin state, one might expect more features to have appeared in the IRMPD spectrum if the quartet had been generated. We note that a shift of the Mo=O stretching frequency of 20 cm$^{-1}$ is predicted by DFT if one assumes oxidation to result in formation of the doublet state, in excellent agreement with the observed shift in the frequency in the experimental spectrum. The quartet state is predicted to lie significantly higher in energy than the doublet, which also argues in favor of assignment of the latter as the spin state. Finally, there is strong evidence to support the doublet state with “folded” geometry that comes from analysis of molecular orbital energies revealed by measurements of vertical and adiabatic electron binding energies by photoelectron spectroscopy. We may therefore conclude that the peak in the IRMPD spectrum of [MoO\(_5\)O(mnt)$_2$]$^{-}$ can be assigned to the Mo=O stretch of the complex in a doublet spin state. The present investigation suggests that the folding of the dithiolene
III and IV, and comparisons of IRMPD and theoretical spectra edged.

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